

Black hole formation in a contracting universe

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Based on work with **Robert Brandenberger** (McGill U.)

JCAP **1611**, 029 (2016) [arXiv:1609.02556]

Λ CDM parameters

Parameter	TT,TE,EE+lowE+lensing+BAO 68% limits
$\Omega_b h^2$	0.02242 ± 0.00014
$\Omega_c h^2$	0.11933 ± 0.00091
$100\theta_{MC}$	1.04101 ± 0.00029
τ	0.0561 ± 0.0071
$\ln(10^{10} A_s)$	3.047 ± 0.014
n_s	0.9665 ± 0.0038

Planck18 [1807.06209]

- $\alpha_s \equiv dn_s/d \ln k = -0.004 \pm 0.013, r < 0.065$ (95 %)
- $f_{NL}^{local} = 0.8 \pm 5.0, f_{NL}^{equil} = -4 \pm 43, f_{NL}^{ortho} = -26 \pm 21$ Planck15 [1502.01592]
- Many inflation models can give you those numbers

But inflation is not the only possibility!

- Ekpyrotic cosmology (contraction with $p = w\rho$, $w \gg 1$) *Khouri et al.*

[hep-th/0103239], Ijjas *et al.* [1404.1265], Lehnert & Wilson-Ewing [1507.08112], Fertig *et al.* [1607.05663]

$$n_s \sim 0.97, \alpha_s \sim \mathcal{O}(-10^{-3}), r \approx 0, |f_{\text{NL}}| \sim \mathcal{O}(1) - \mathcal{O}(10)$$

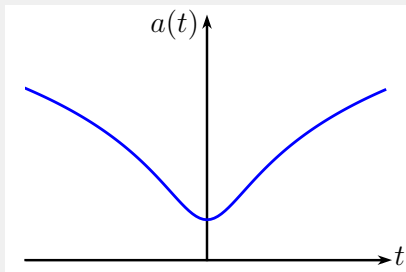
- String gas cosmology (quasi-static, thermal) *Brandenberger & Vafa [89], Chen et al.*

[0712.2477], Brandenberger *et al.* [1403.4927], Brandenberger [1505.02381]

$$n_s < 1, r \ll 1, |f_{\text{NL}}| \ll 1, n_t \approx 1 - n_s > 0$$

- And other scenarios *See review: Brandenberger & Peter [1603.05834]*

Bouncing Cosmology



- Many alternatives are ‘**bouncing**’ scenarios: the Big Bang is replaced by a bounce, before which the Universe was contracting
- Can we find generic **predictions** or imprints of a **contracting Universe** before the Big Bang/bounce?
- Take a minimal assumption: consider a contracting universe with matter as we know it today (i.e., dust, radiation, etc.)

Hydrodynamical fluid in a contracting universe

- Energy density and pressure:

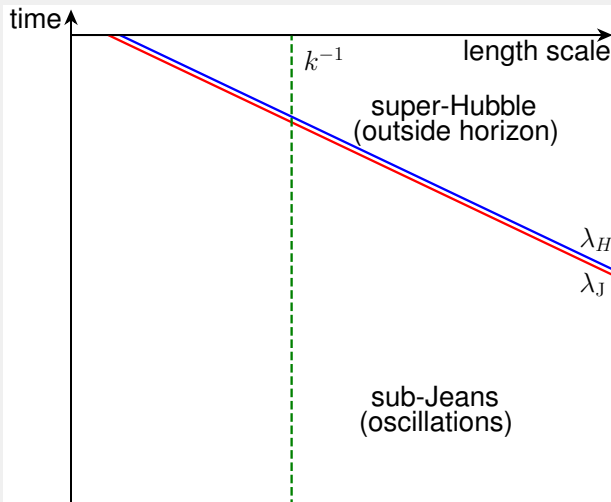
$$\rho(t, \mathbf{x}) = \bar{\rho}(t) + \delta\rho(t, \mathbf{x}), \quad p(t, \mathbf{x}) = \bar{p}(t) + \delta p(t, \mathbf{x})$$

$$\bar{p} = w\bar{\rho}, \quad \delta p = c_s \delta\rho \text{ (constant entropy)}, \quad c = 1$$

- Dust: $w = c_s^2 \ll 1$; radiation $w = c_s^2 = 1/3$
- Density contrast in Fourier space: $\delta_k \equiv \delta\rho_k/\bar{\rho}$
- Jeans scale: $k_J \equiv \sqrt{4\pi G_N \bar{\rho}/c_s}$, $\lambda_J \sim 1/k_J$
- Sub-Jeans scales ($\lambda < \lambda_J$) \implies oscillations ($\delta_k(t) \sim e^{\pm i c_s k t}$)
- Super-Jeans scales ($\lambda > \lambda_J$) \implies gravitational instability ($\delta_k(t) \nearrow$)
- Can this lead to black holes?

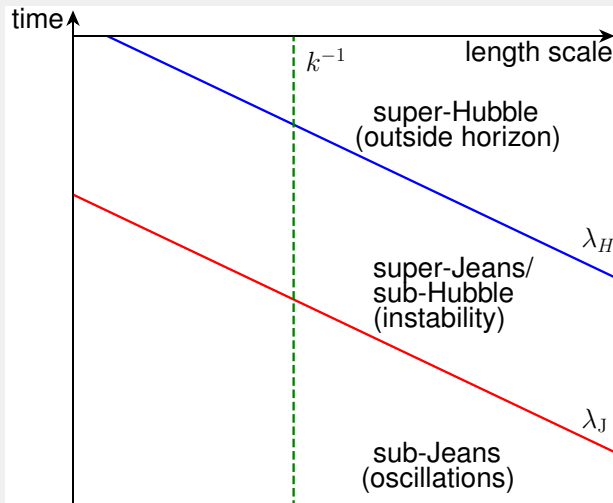
Jeans and Hubble scales

- Jeans radius: $\lambda_J = c_s(a|H|)^{-1}$; Hubble radius: $\lambda_H = (a|H|)^{-1}$
- Radiation $\implies c_s^2 = 1/3 \implies \lambda_J \sim \lambda_H$



Jeans and Hubble scales

- Jeans radius: $\lambda_J = c_s(a|H|)^{-1}$; Hubble radius: $\lambda_H = (a|H|)^{-1}$
- Dust $\implies c_s^2 \ll 1 \implies \lambda_J \ll \lambda_H$



Black hole formation probability

- Specify initial conditions (e.g., quantum vacuum, thermal state, etc.)
→ get full solution for $\delta_k(t)$ by solving the linearized Einstein equations
- Probability of black hole formation (Press-Schechter formalism):

$$\text{Prob}(R, t) = \sqrt{\frac{2}{\pi}} \frac{1}{\sigma} \int_{\delta_c}^{\infty} d\delta e^{-\delta^2/2\sigma^2},$$

$$\sigma^2(R, t) = \frac{1}{2\pi^2} \int_0^{\infty} dk k^2 W^2(kR) |\delta_k(t)|^2,$$

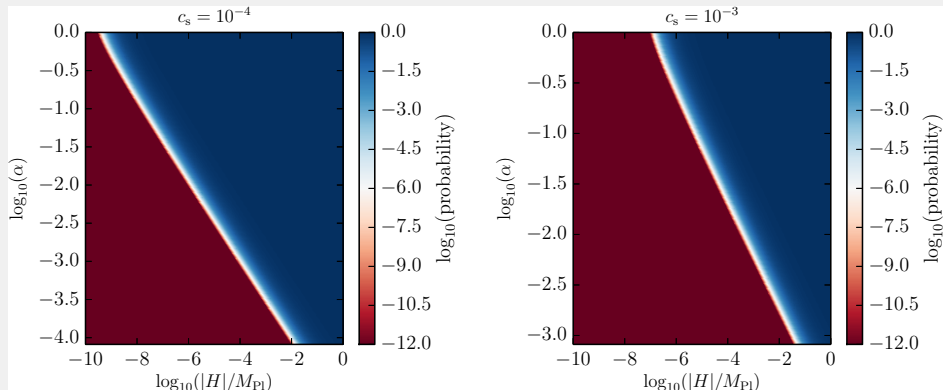
$W(kR)$ = Window function

- High probability when $\sigma \gtrsim \delta_c$

Black hole formation probability

- Starting with a quantum vacuum [JQ & Brandenberger \[1609.02556\]](#)

$$\alpha \equiv R_{\text{BH}}/|H|^{-1}$$

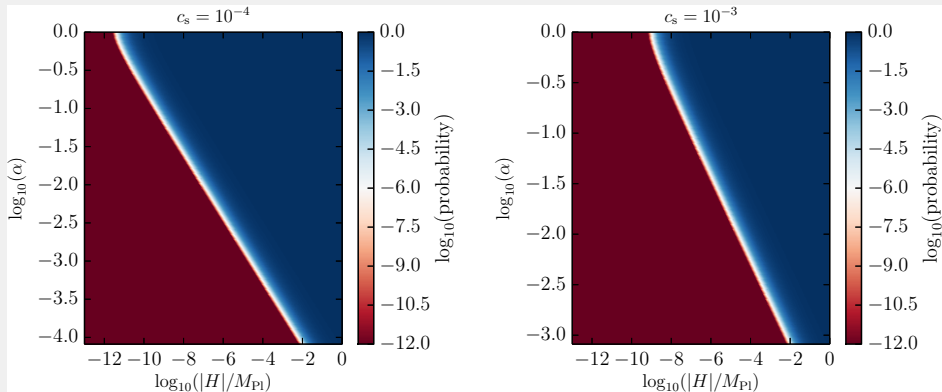


dust, super-Jeans/sub-Hubble: $|\delta_k(t)|^2 \sim k^{-3} c_s^{-5} H(t)^2$

Black hole formation probability

- Starting with a thermal state at $|H_{\text{ini}}| = 10^{-16} M_{\text{Pl}}$ JQ & Brandenberger [1609.02556]

$$\alpha \equiv R_{\text{BH}}/|H|^{-1}$$



Implications and discussion

- If a long-enough phase of dust-dominated contraction lasts, Hubble-size black holes form well before Planckian densities:

$$R_{\text{BH}} \sim c_s^{-5/2} \ell_{\text{Pl}} \text{ (quantum)}; R_{\text{BH}} \sim c_s^{-18/5} \ell_{\text{Pl}} \text{ (thermal)} \quad \text{JQ \& Brandenberger}$$

[1609.02556]

- E.g. (quantum initial conditions):

$$c_s = 10^{-10} \implies R_{\text{BH}} \sim 10^{25} \ell_{\text{Pl}} \sim 10^{-10} \text{ m},$$

$$M_{\text{BH}} \sim 10^{20} \text{ g} \sim 10^{-13} M_{\odot} \longrightarrow \text{primordial black hole?}$$

- E.g. 2: $c_s = 10^{-5} \implies R_{\text{BH}} \sim 10^{12} \ell_{\text{Pl}} \sim 10^{-23} \text{ m}$, $M_{\text{BH}} \sim 10^7 \text{ g} \longrightarrow$ evaporated black hole remnant?

- Harder to form black holes if c_s increases (e.g., no radiation black holes before Planckian densities)
- But easier to form black holes if there are structures already in the Universe (rather than just a quantum vacuum)

Outlook

- First estimate of black hole formation in a contracting universe → analysis could be refined
- Can black holes pass through a bounce? Carr *et al.* [1104.3796, 1402.1437, 1701.05750, 1704.02919]
- If so, one can use observations to constrain bouncing cosmological models Chen *et al.* [1609.02571]
- Could there be specific gravitational wave signals? Barrau *et al.* [1711.05301]
- Also, could black holes play a role in the bounce itself? E.g., black holes at the **string scale** Mathur [0803.3727], Masoumi [1505.06787], JQ *et al.* [1809.01658]

Thank you for your attention!

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